

GENERAL TAX STRUCTURES AND THE LÉVY INSURANCE RISK MODEL.

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Abstract

In the spirit of [2, 3] we consider a Lévy insurance risk model with tax payments of a more general structure than in the aforementioned papers that was also considered in [1]. In terms of scale functions, we establish three fundamental identities of interest which have stimulated a large volume of actuarial research in recent years. That is to say, the two sided exit problem, the net present value of tax paid until ruin as well as a generalized version of the Gerber-Shiu function. The method we appeal to differs from [2, 3] in that we appeal predominantly to excursion theory.

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1. Introduction and main results

Recent advances in the analysis of the ubiquitous ruin problem from the theory of insurance risk has seen a tendency to replace the classical Cramér-Lundberg surplus process with a general spectrally negative Lévy processes; see for example [3, 12, 13, 24] to name but a few. In that case the surplus process is commonly referred to as a *Lévy insurance risk process*. Although moving to this more complex setting, arguably, does not bring any more realistic features to the table than are already on offer in the classical Cramér-Lundberg model, a clear mathematical advantage has emerged. Working with Lévy insurance risk processes forces one to approach the problem of ruin via excursion or fluctuation theory which does not use specific features of the underlying Lévy process other than a generic path decomposition of the process in terms of excursions from its maximum which manifests itself in the form of a Poisson point process.

In this paper, we continue in this vein and build on ideas of Lévy insurance risk processes with tax which were introduced and studied in [2, 3, 1]. Specifically we introduce a more general tax structure and therewith we establish, for the aggregate surplus process, new identities for the two sided exit problem, a generalized version of the Gerber-Shiu function as well the net present value of tax paid until ruin.

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Henceforth the process $X = \{X_t : t \geq 0\}$ with probabilities $\{\mathbb{P}_x : x \in \mathbb{R}\}$ and natural filtration $\{\mathcal{F}_t : t \geq 0\}$ will denote a spectrally negative Lévy process with the usual exclusion of processes in the latter class which have monotone paths (that is to say a pure increasing linear drift and the negative of a subordinator). For convenience we shall always denote \mathbb{P}_0 by \mathbb{P} . Let

$$\psi(\theta) = \log \mathbb{E}(e^{\theta X_1})$$

be the Laplace exponent of X which is known to be finite for at least $\theta \in [0, \infty)$ in which case it is a strictly convex and infinitely differentiable function. The asymptotic behaviour of X is characterized by $\psi'(0+)$, so that X drifts to $\pm\infty$ (oscillates) accordingly as $\pm\psi'(0+) > 0$ ($\psi'(0+) = 0$). When X plays the role of the surplus process, it is usual to make the assumption that $\psi'(0+) > 0$ which is equivalent to the *net profit condition* in the case that X is a Cramér-Lundberg process. However, this condition is not necessary for any of the forthcoming analysis.

Denote by $S = \{S_t : t \geq 0\}$ the process which describes the running supremum of X , that is to say, $S_t = \sup_{s \leq t} X_s$ for each $t \geq 0$. Following [1], we are interested in modeling tax payments from the Lévy insurance process in such a way that the cumulative payment until time t is given by

$$\int_0^t \gamma(S_u) dS_u.$$

where $\gamma : [0, \infty) \rightarrow [0, 1)$ is a measurable function which satisfies

$$\int_0^\infty (1 - \gamma(s)) ds = \infty. \quad (1)$$

In that case the aggregate surplus process, the primary object of our study, is given by

$$U_t := X_t - \int_0^t \gamma(S_u) dS_u. \quad (2)$$

In the special case that γ is a constant in $(0, 1)$ our Lévy insurance risk process with tax agrees with the model introduced in [2, 3]. In the case that $\gamma = 0$, we are back to a regular Lévy insurance risk process. Note also that processes of the form (2) constitute a subclass of controlled Lévy risk processes, the latter being of popular interest in recent literature; see for example [5, 16, 17].

In order to state the results alluded to above which concern path functionals of U , we must first introduce more notation. As is now usual when studying Lévy risk processes, a key element of the analysis involves the use of scale functions, defined as follows. For every $q \geq 0$ there exists a function $W^{(q)} : \mathbb{R} \rightarrow [0, \infty)$ such that $W^{(q)}(x) = 0$ for all $x < 0$ and otherwise is almost everywhere differentiable on $[0, \infty)$ satisfying,

$$\int_0^\infty e^{-\lambda x} W^{(q)}(x) dx = \frac{1}{\psi(\lambda) - q}, \quad \text{for } \lambda > \Phi(q), \quad (3)$$

where $\Phi(q)$ is the largest solution to the equation $\psi(\theta) = q$ (there are at most two). We shall write for short $W^{(0)} = W$. It is known that when X has paths of unbounded

variation, the scale functions $W^{(q)}$ are continuously differentiable on $(0, \infty)$, and when X has paths of bounded variation, they are almost everywhere differentiable. In either case we shall denote by $W^{(q)'}$ the associated density. It is also known that if X has a Gaussian component then $W^{(q)}$ is twice continuously differentiable on $(0, \infty)$. In the case of bounded variation paths, there is a discontinuity in $W^{(q)}$ at the origin which does not depend on q and in the case of unbounded variation paths, there is no discontinuity in $W^{(q)}$ at the origin for all $q \geq 0$.

There exists a well known exponential change of measure that one may perform for spectrally negative Lévy processes,

$$\frac{d\mathbb{P}_x^\vartheta}{d\mathbb{P}_x} \Big|_{\mathcal{F}_t} = e^{\vartheta(X_t - x) - \psi(\vartheta)t} \quad (4)$$

for $x \in \mathbb{R}$ and $\vartheta \geq 0$ under which X remains within the class of spectrally negative Lévy processes. In particular if $\nu(dx)$ is the Lévy measure of $-X$ under \mathbb{P} then $e^{-\vartheta x} \nu(dx)$ is its Lévy measure under \mathbb{P}^ϑ . It will turn out to be useful to introduce an additional parameter to the scale functions described above in the light of this change of measure. Henceforth we shall refer to the functions W_ϑ where $\vartheta \geq 0$ as the functions that play the role of the scale functions defined in the previous paragraph but when considered under the measures \mathbb{P}^ϑ .

Next define

$$\tau_a^+ := \inf\{t > 0 : U_t > a\} \text{ and } \tau_0^- := \inf\{t > 0 : U_t < 0\}$$

with the convention $\inf \emptyset = \infty$. For $s \geq x$ define

$$\bar{\gamma}(s) := s - \int_x^s \gamma(y) dy = x + \int_x^s (1 - \gamma(y)) dy. \quad (5)$$

By differentiating (5) latter we note that since $\gamma \in [0, 1]$ it follows that $\bar{\gamma}$ is a strictly increasing function. Moreover, since it is continuous it has a well defined inverse on $[x, \infty)$ which we denote by $\bar{\gamma}^{-1}$.

We may now present the three main results of this paper as promised. Their proofs will be given in the subsequent sections.

For each $x > 0$, the process $L_t := S_t - x$, $t \geq 0$, serves as a local time at 0 for the Markov process $Y := S - X$ under \mathbb{P}_x . Write $L^{-1} := \{L_t^{-1} : t \geq 0\}$ for the right continuous inverse of L .

Theorem 1.1. (Two sided exit problem.) *For any $x > 0$ and $q \geq 0$, we have*

$$\mathbb{E}_x \left[e^{-q\tau_a^+} \mathbb{I}_{\{\tau_a^+ < \tau_0^-\}} \right] = \exp \left\{ - \int_x^{a \wedge x} \frac{W^{(q)'}(y)}{W^{(q)}(y)(1 - \gamma(\bar{\gamma}^{-1}(y)))} dy \right\}. \quad (6)$$

Theorem 1.2. (Net present value of tax paid until ruin.) *For any $x > 0$ and $q \geq 0$, we have*

$$\mathbb{E}_x \left[\int_0^{\tau_0^-} e^{-qu} \gamma(S_u) dS_u \right] = \int_x^\infty \exp \left\{ - \int_x^t \frac{W^{(q)'(\bar{\gamma}(s))}}{W^{(q)}(\bar{\gamma}(s))} ds \right\} \gamma(t) dt. \quad (7)$$

Theorem 1.3. *For each $t \geq 0$ let $S_t^U := \sup_{s \leq t} U_s$ and let $\kappa = L_{L_{\tau_0^-}^-}^{-1}$, the last moment that tax is paid before ruin. Denote by ν the Lévy measure of $-X$. For any $z > 0, y \geq 0, \theta \geq y$ and $\alpha, \beta \geq 0$, we have*

$$\begin{aligned} & \mathbb{E}_x \left(e^{-\alpha\kappa - \beta(\tau_0^- - \kappa)}; S_{\tau_0^-}^U \in d\theta, U_{\tau_0^-} \in dy, -U_{\tau_0^-} \in dz \right) \\ &= \frac{1}{1 - \gamma(\bar{\gamma}^{-1}(\theta))} \exp \left\{ - \int_x^\theta \frac{W^{(\alpha)'(v)}}{W^{(\alpha)}(v)(1 - \gamma(\bar{\gamma}^{-1}(v)))} dv \right\} \\ & \quad \cdot \left[\left\{ W^{(\beta)'(\theta - y)} - \frac{W^{(\beta)'(\theta)}}{W^{(\beta)}(\theta)} W^{(\beta)}(\theta - y) \right\} \nu(y + dz) \mathbf{1}_{\{y < \theta\}} dy \right. \\ & \quad \left. + W^{(\beta)}(0+) \nu(\theta + dz) \delta_\theta(dy) \right] d\theta \quad (8) \end{aligned}$$

where $\delta_\theta(dy)$ is the Dirac measure which assigns unit mass to the point θ . Furthermore, we also have,

$$\begin{aligned} & \mathbb{E}_x \left(e^{-\alpha\kappa - \beta(\tau_0^- - \kappa)}; S_{\tau_0^-}^U \in d\theta, U_{\tau_0^-} = 0 \right) \\ &= \frac{1}{1 - \gamma(\bar{\gamma}^{-1}(\theta))} \exp \left\{ - \int_x^\theta \frac{W^{(\alpha)'(y)}}{W^{(\alpha)}(y)(1 - \gamma(\bar{\gamma}^{-1}(y)))} dy \right\} \\ & \quad \cdot \frac{\sigma^2}{2} \left\{ \frac{W^{(\beta)'(\theta)^2}}{W^{(\beta)}(\theta)} - W^{(\beta)''(\theta)} \right\} d\theta, \quad (9) \end{aligned}$$

where σ is the Gaussian coefficient in the Lévy-Itô decomposition of X .

Remark 1. When $\gamma \in (0, 1)$ is a constant, we note that the expressions (6) and (7) agree with formulas (3.1) and (3.2) in [3]. Indeed we have $\bar{\gamma}(s) = s(1 - \gamma) + \gamma x$. For Theorem 1.1 we have

$$\mathbb{E}_x \left[e^{-q\tau_a^+} \mathbb{1}_{\{\tau_a^+ < \tau_0^-\}} \right] = \exp \left\{ - \int_x^a \frac{W^{(q)'(y)}}{W^{(q)}(y)(1 - \gamma)} dy \right\} = \left(\frac{W^{(q)}(x)}{W^{(q)}(a)} \right)^{1/(1-\gamma)}.$$

For Theorem 1.2 we have by two changes of variables

$$\begin{aligned} \int_x^\infty \exp \left\{ - \int_x^t \frac{W^{(q)'(\bar{\gamma}(s))}}{W^{(q)}(\bar{\gamma}(s))} ds \right\} \gamma(t) dt &= \gamma \int_x^\infty \exp \left\{ - \frac{1}{1 - \gamma} \int_x^{\bar{\gamma}(t)} \frac{W^{(q)'(y)}}{W^{(q)}(y)} dy \right\} dt \\ &= \gamma \int_x^\infty \left(\frac{W^{(q)}(x)}{W^{(q)}(\bar{\gamma}(t))} \right)^{1/(1-\gamma)} dt \\ &= \frac{\gamma}{1 - \gamma} \int_x^\infty \left(\frac{W^{(q)}(x)}{W^{(q)}(u)} \right)^{1/(1-\gamma)} du, \end{aligned}$$

which is formula (3.2) in [3].

Theorem 1.3 on the other hand gives a new result for the setting of [3]. In particular,

we have for $y, z > 0$ and $y \leq \theta$,

$$\begin{aligned} & \mathbb{E}_x \left(e^{-\alpha\kappa - \beta(\tau_0^- - \kappa)}; S_{\tau_0^-}^U \in d\theta, U_{\tau_0^-} \in dy, -U_{\tau_0^-} \in dz \right) \\ &= \frac{1}{1-\gamma} \left(\frac{W^{(\alpha)}(x)}{W^{(\alpha)}(\theta)} \right)^{1/(1-\gamma)} \times \\ & \left[\left\{ W^{(\beta)'}(\theta - y) - \frac{W^{(\beta)'}(\theta)}{W^{(\beta)}(\theta)} W^{(\beta)}(\theta - y) \right\} \nu(y + dz) \mathbf{1}_{\{y < \theta\}} dy + W^{(\beta)}(0+) \nu(\theta + dz) \delta_\theta(dy) \right] d\theta \end{aligned}$$

and

$$\begin{aligned} & \mathbb{E} \left(e^{-\alpha\kappa - \beta(\tau_0^- - \kappa)}; S_{\tau_0^-}^U \in d\theta, U_{\tau_0^-} = 0 \right) \\ &= \frac{\sigma^2}{2(1-\gamma)} \left(\frac{W^{(\alpha)}(x)}{W^{(\alpha)}(\theta)} \right)^{1/(1-\gamma)} \left\{ \frac{W^{(\beta)'}(\theta)^2}{W^{(\beta)}(\theta)} - W^{(\beta)''}(\theta) \right\} d\theta. \end{aligned}$$

Finally note that when $\gamma = 0$ and the process U agrees with the Lévy insurance risk process X , the last two formulae above give us two new expressions for the time value of the overall maximal wealth accumulated prior to ruin, the wealth immediately before ruin and the deficit at ruin.

Remark 2. Following a preprint of this article, further identities concerning the process U have been treated by Renaud [23] using a different technique.

Remark 3. One major criticism of working with scale functions is that, in principle, one has only solved the problems of interest up to inverting the Laplace transform (3). However, in the last year there have been a number of developments in the theory of scale functions which has seen a large number of explicit examples appearing in the literature; including the case of Cramér-Lundberg models. See for example [9, 11, 18, 20]. The paper [25] also gives recipes for evaluating scale functions numerically.

2. Proofs of Main results

We begin this section by pointing out some important features of the running supremum of the aggregate process (2) which turns out to be key in our use of excursion theory in the forthcoming proofs.

Lemma 2.1. *We have that*

$$S_t^U = S_t - \int_0^t \gamma(S_s) dS_s \quad (10)$$

and that the random times $\{t \geq 0 : U_t = S_t^U\}$ agree precisely with $\{t \geq 0 : X_t = S_t\}$.

Proof. Note that on the one hand

$$S_t^U = \sup_{s \leq t} \left\{ X_s - \int_0^s \gamma(S_u) dS_u \right\} \geq \sup_{s \leq t} X_s - \int_0^t \gamma(S_u) dS_u = S_t - \int_0^t \gamma(S_u) dS_u. \quad (11)$$

On the other hand since $X_t \leq S_t$ we have

$$U_t \leq S_t - \int_0^t \gamma(S_u) dS_u = \int_0^t (1 - \gamma(S_u)) dS_u + x$$

and hence, since $\gamma(y) \in [0, 1]$ for all $y \geq 0$,

$$S_t^U \leq \sup_{s \leq t} \int_0^s (1 - \gamma(S_u)) dS_u + x = \int_0^t (1 - \gamma(S_u)) dS_u + x = S_t - \int_0^t \gamma(S_u) dS_u \quad (12)$$

Together (12) and (11) imply (10). Now suppose that $t' \in \{t \geq 0 : X_t = S_t\}$. This implies that

$$U_{t'} = X_{t'} - \int_0^{t'} \gamma(S_u) dS_u = S_{t'} - \int_0^{t'} \gamma(S_u) dS_u = S_{t'}^U$$

and hence $t' \in \{t \geq 0 : U_t = S_t^U\}$. On the other hand, if $t'' \in \{t \geq 0 : U_t = S_t^U\}$, then

$$X_{t''} - \int_0^{t''} \gamma(S_u) dS_u = U_{t''} = S_{t''}^U = S_{t''} - \int_0^{t''} \gamma(S_u) dS_u$$

showing that $X_{t''} = S_{t''}$ and hence $t'' \in \{t \geq 0 : X_t = S_t\}$. \square

For the remaining proofs we shall also make heavy use of excursion theory for the process $S - X$ for which we refer to [6] for background reading. We shall spend a moment here setting up some necessary notation which will be used throughout the remainder of the paper. The Poisson point process of excursions indexed by local time shall be denoted by $\{(t, \epsilon_t) : t \geq 0\}$ where

$$\epsilon_t = \{\epsilon_t(s) := X_{L_t^{-1}} - X_{L_{t-}^{-1} + s} : 0 < s \leq L_t^{-1} - L_{t-}^{-1}\}$$

whenever $\sigma(\epsilon_t) := L_t^{-1} - L_{t-}^{-1} > 0$. Accordingly we refer to a generic excursion as $\epsilon(\cdot)$ (or just ϵ for short as appropriate) belonging to the space \mathcal{E} of canonical excursions. The intensity measure of the process $\{(t, \epsilon_t) : t \geq 0\}$ is given by $dt \times dn$ where n is a measure on the space of excursions (the excursion measure). An n -measurable functional of the canonical excursion which will be of prime interest is $\bar{\epsilon} = \sup_{s \geq 0} \epsilon(s)$. A useful formula for this functional that we shall make use of is the following (cf. [14])

$$n(\bar{\epsilon} > x) = \frac{W'(x)}{W(x)} \quad (13)$$

providing that x is not a point of discontinuity in the derivative of W (which is only a concern when X has paths of bounded variation, in which case there are at most a countable number).

Lemma 2.1 also has an important bearing on the process of excursions described above. Indeed, from the identity (10) we note that $L_t = s$, or equivalently $S_t = x + s$, under \mathbb{P}_x if and only if $S_t^U = \bar{\gamma}(x + s)$. It follows that the event that U climbs from x to height $\bar{\gamma}(x + s)$ for the first time coincides with the event that process X climbing from x to $x + s$ for the first time. Consequently, $L_s^{-1} = \tau_{\bar{\gamma}(x+s)}^+$, or equivalently $\tau_a^+ = L_{\bar{\gamma}^{-1}(a)-x}^{-1}$, under \mathbb{P}_x . Moreover, the excursions of U away from its maximum agree precisely with $\{(t, \epsilon_t) : t \geq 0\}$.

Proof of Theorem 1.1. Taking account of the remarks following Lemma 2.1 we have that the event $\{\tau_a^+ < \tau_0^-\}$ is the same as

$$\{\bar{\epsilon}_s \leq \bar{\gamma}(x + s), \forall 0 \leq s < \bar{\gamma}^{-1}(a) - x\}.$$

Then, for $x > 0$,

$$\begin{aligned}
\mathbb{P}_x(\tau_a^+ < \tau_0^-) &= \mathbb{P}_x(\bar{\epsilon}_s \leq \bar{\gamma}(x+s), \forall 0 \leq s < \bar{\gamma}^{-1}(a) - x) \\
&= \exp \left\{ - \int_0^{\bar{\gamma}^{-1}(a)-x} n(\bar{\epsilon} > \bar{\gamma}(x+s)) ds \right\} \\
&= \exp \left\{ - \int_0^{\bar{\gamma}^{-1}(a)-x} \frac{W'(\bar{\gamma}(x+s))}{W(\bar{\gamma}(x+s))} ds \right\} \\
&= \exp \left\{ - \int_x^a \frac{W'(y)}{W(y)(1 - \gamma(\bar{\gamma}^{-1}(y)))} dy \right\},
\end{aligned}$$

where we change the variable making use of the fact that, since $\bar{\gamma}(\bar{\gamma}^{-1}(s)) = s$, we have from the chain rule

$$\frac{d}{ds} \bar{\gamma}^{-1}(s) = \frac{1}{\bar{\gamma}'(\bar{\gamma}^{-1}(s))} = \frac{1}{1 - \gamma(\bar{\gamma}^{-1}(s))}.$$

Next, note that

$$\begin{aligned}
\mathbb{P}_x^{\Phi(q)}(\tau_a^+ < \tau_0^-) &= \mathbb{E}_x \left[e^{\Phi(q)(X_{\tau_a^+} - x) - q\tau_a^+} \mathbb{I}_{\{\tau_a^+ < \tau_0^-\}} \right] \\
&= \exp \{ \Phi(q) (\bar{\gamma}^{-1}(a) - x) \} \mathbb{E}_x \left[e^{-q\tau_a^+} \mathbb{I}_{\{\tau_a^+ < \tau_0^-\}} \right] \quad (14)
\end{aligned}$$

where we have appealed to the change of measure (4) with $\vartheta = \Phi(q)$ and the final equality follows by virtue of the fact that on $\{\tau_a^+ < \infty\}$

$$\begin{aligned}
X_{\tau_a^+} &= U_{\tau_a^+} + \int_0^{\tau_a^+} \gamma(S_u) dS_u \\
&= a + \int_0^{L_{\bar{\gamma}^{-1}(a)-x}} \gamma(S_u) dS_u \\
&= a + \int_x^{\bar{\gamma}^{-1}(a)} \gamma(y) dy \\
&= \bar{\gamma}^{-1}(a)
\end{aligned}$$

where in the third equality we have made the change of variable $y = S_u^{-1}$. Note also that it is known (cf. Chapter 8 of [14]) that for $q, x \geq 0$,

$$W^{(q)}(x) = e^{\Phi(q)x} W_{\Phi(q)}(x) \quad (15)$$

and hence

$$\frac{W'_{\Phi(q)}(x)}{W_{\Phi(q)}(x)} = \frac{W^{(q)'}(x)}{W^{(q)}(x)} - \Phi(q). \quad (16)$$

Piecing together (14), (15) and (16) we get

$$\begin{aligned}
& \mathbb{E}_x \left[e^{-q\tau_a^+} \mathbb{I}_{\{\tau_a^+ < \tau_0^-\}} \right] \\
&= \mathbb{P}_x^{\Phi(q)}(\tau_a^+ < \tau_0^-) \exp \left\{ -\Phi(q) (\bar{\gamma}^{-1}(a) - x) \right\} \\
&= \exp \left\{ -\int_x^a \frac{W_{\Phi(q)}'(y)}{W_{\Phi(q)}(y)(1 - \gamma(\bar{\gamma}^{-1}(y)))} dy \right\} \exp \left\{ -\Phi(q) (\bar{\gamma}^{-1}(a) - x) \right\} \\
&= \exp \left\{ -\int_x^a \frac{W^{(q)'}(y)}{W^{(q)}(y)(1 - \gamma(\bar{\gamma}^{-1}(y)))} dy \right\},
\end{aligned}$$

where we have also used the fact that

$$\int_x^a \frac{1}{1 - \gamma(\bar{\gamma}^{-1}(y))} dy = \bar{\gamma}^{-1}(a) - \bar{\gamma}^{-1}(x) = \bar{\gamma}^{-1}(a) - x.$$

The proof is now complete. \square

Proof of Theorem 1.2. The proof builds on the experience of the calculations in the previous proof. We note that the process S does not increase on the time interval $(L_{L_{\tau_0^-}^{-1}}, \tau_0^-)$ and hence

$$\begin{aligned}
& \mathbb{E}_x \left[\int_0^{\tau_0^-} e^{-qu} \gamma(S_u) dS_u \right] \\
&= \mathbb{E}_x \left[\int_0^{L_{L_{\tau_0^-}^{-1}}} e^{-qu} \gamma(L_u + x) dL_u \right] \\
&= \mathbb{E}_x \left[\int_0^\infty \mathbb{I}_{\{t < L_{\tau_0^-}\}} e^{-qL_t^{-1}} \gamma(t + x) dt \right] \\
&= \int_0^\infty \mathbb{E}_x \left[e^{-qL_t^{-1}} \mathbb{I}_{\{\bar{\epsilon}_s \leq \bar{\gamma}(x+s), \forall 0 \leq s \leq t\}} \right] \gamma(t + x) dt \\
&= \int_0^\infty e^{-\Phi(q)t} \mathbb{P}_x^{\Phi(q)}(\bar{\epsilon}_s \leq \bar{\gamma}(x+s), \forall 0 \leq s \leq t) \gamma(t + x) dt \\
&= \int_0^\infty e^{-\Phi(q)t} \exp \left\{ -\int_0^t n_{\Phi(q)}(\bar{\epsilon} > \bar{\gamma}(x+s)) ds \right\} \gamma(t + x) dt \\
&= \int_0^\infty e^{-\Phi(q)t} \exp \left\{ -\int_0^t \frac{W'_{\Phi(q)}(\bar{\gamma}(x+s))}{W_{\Phi(q)}(\bar{\gamma}(x+s))} ds \right\} \gamma(t + x) dt \\
&= \int_0^\infty \exp \left\{ -\int_0^t \frac{W^{(q)' }(\bar{\gamma}(x+s))}{W^{(q)}(\bar{\gamma}(x+s))} ds \right\} \gamma(t + x) dt
\end{aligned}$$

where in the fifth equality the measure $n_{\Phi(q)}$ plays the role of n under $\mathbb{P}^{\Phi(q)}$, in the penultimate equality we have used (13) and the final equality uses (16). The proof is completed by applying a straightforward change of variables. \square

Before turning to the proof of Theorem 1.3, we need first to prove an additional auxiliary result. To this end, define $\rho_a = \inf\{s > 0 : \epsilon(s) > a\}$, the first passage

time above a of the canonical excursion ϵ . We also need the first passage times for the underlying Lévy process X ,

$$T_x^+ = \inf\{t > 0 : X_t > x\} \text{ and } T_x^- = \inf\{t > 0 : X_t < x\}$$

for all $x \in \mathbb{R}$.

Lemma 2.2. *For any $z > 0$, $y \in [0, a]$ and $q \geq 0$, we have*

$$\begin{aligned} & n(e^{-q\rho_a}; a - \epsilon(\rho_{a-}) \in dy, \epsilon(\rho_a) - a \in dz) \\ &= \left\{ W^{(q)'}(a-y) - \frac{W^{(q)'(a)}}{W^{(q)}(a)} W^{(q)}(a-y) \right\} \nu(y+dz) \mathbf{1}_{\{y < a\}} dy + W^{(q)}(0+) \nu(a+dz) \delta_a(dy) \end{aligned}$$

where $\delta_a(dy)$ is the Dirac measure assigning unit mass to the point a and

$$n(e^{-q\rho_a}; \epsilon(\rho_a) = a) = \frac{\sigma^2}{2} \left\{ \frac{W^{(q)'(a)^2}{W^{(q)}(a)} - W^{(q)''(a)} \right\}.$$

Proof. Recall that $Y = S - X$. For the latter process introduce its first passage time

$$\varsigma_a = \inf\{t > 0 : Y_t > a\}.$$

By a classical application of the compensation formula (see for example the treatment of a related problem in [4]) we have for $z > 0$, $y \in [0, a]$ and $q \geq 0$ that

$$\begin{aligned} & \mathbb{E}(e^{-q\varsigma_a}; a - Y_{\varsigma_a-} \in dy, Y_{\varsigma_a} - a \in dz) \\ &= \mathbb{E} \left[\sum_{t \geq 0} e^{-q(L_t^{-1} + \rho_a(\epsilon_t))} \mathbb{I}_{\{\sup_{s < t} \bar{\epsilon}_s \leq a, \rho_a(\epsilon_t) < \sigma(\epsilon_t), a - \epsilon_t(\rho_{a-}) \in dy, \epsilon_t(\rho_a) - a \in dz\}} \right] \\ &= \int_0^\infty \mathbb{E} \left[e^{-qL_t^{-1}} \mathbb{I}_{\{\sup_{s \leq t} \bar{\epsilon}_s \leq a\}} \right] dt \cdot n(e^{-q\rho_a}; a - \epsilon(\rho_{a-}) \in dy, \epsilon(\rho_a) - a \in dz) \\ &= \int_0^\infty e^{-\Phi(q)t} e^{-n_{\Phi(q)}(\bar{\epsilon} > a)t} dt \cdot n(e^{-q\rho_a}; a - \epsilon(\rho_{a-}) \in dy, \epsilon(\rho_a) - a \in dz) \\ &= \int_0^\infty e^{-\Phi(q)t} \exp \left\{ -\frac{W'_{\Phi(q)}(a)}{W_{\Phi(q)}(a)} t \right\} dt \cdot n(e^{-q\rho_a}; a - \epsilon(\rho_{a-}) \in dy, \epsilon(\rho_a) - a \in dz) \\ &= \int_0^\infty \exp \left\{ -\frac{W^{(q)'(a)}}{W^{(q)}(a)} t \right\} dt \cdot n(e^{-q\rho_a}; a - \epsilon(\rho_{a-}) \in dy, \epsilon(\rho_a) - a \in dz) \quad (17) \end{aligned}$$

where in the first equality the time index runs over local times and the sum is the usual shorthand for integration with respect to the Poisson counting measure of excursions, and for the second equality we need the quasi left continuity for subordinator L^{-1} .

On the other hand, according to Theorem 1 of [22], we have that

$$\begin{aligned} & \mathbb{E}(e^{-q\varsigma_a}; a - Y_{\varsigma_a-} \in dy, Y_{\varsigma_a} - a \in dz) \\ &= \int_0^\infty \exp \left\{ -\frac{W^{(q)'(a)}}{W^{(q)}(a)} t \right\} dt \\ & \quad \times \left(W^{(q)'(a-y)} - \frac{W^{(q)'(a)}}{W^{(q)}(a)} W^{(q)}(a-y) \right) \nu(y+dz) dy \quad (18) \end{aligned}$$

By comparing the left and right hand sides of (17) and (18) we thus have that

$$\begin{aligned} & n(e^{-q\rho_a}; a - \epsilon(\rho_{a-}) \in dy, \epsilon(\rho_a) - a \in dz) \\ &= \left(W^{(q)'}(a-y) - \frac{W^{(q)'(a)}}{W^{(q)}(a)} W^{(q)}(a-y) \right) \nu(y+dz)dy \end{aligned}$$

as claimed.

In order to write down an identity for $n(e^{-q\rho_a}; a - \epsilon(\rho_{a-}) \in dy, \epsilon(\rho_a) - a \in dz)$ in the case that $y = a$ and $z > 0$, note that this can only occur when the excursion immediately jumps from zero at time zero over the level a . This is not possible when X has paths of unbounded variation; see for example [19]. However, when X has paths of bounded variation, it is known (cf. [21] for example) that excursions begin with a jump and the entrance law is given by $d^{-1}\nu(dx)$. Here $d > 0$ is the drift term of X when it is uniquely written in the form $X_t = dt - \xi_t$ where ξ is a pure jump subordinator. It is known (cf. Lemma 8.6 of [14]) that, when X has paths of bounded variation, we may write $d^{-1} = W^{(q)}(0+)$ for all $q \geq 0$. Moreover, when X has paths of unbounded variation, we have $W^{(q)}(0+) = 0$ for all $q \geq 0$. It follows that we may now generically write for $z > 0$,

$$\begin{aligned} n(e^{-q\rho_a}; a - \epsilon(\rho_{a-}) \in \{a\}, \epsilon(\rho_a) - a \in dz) &= n(\epsilon(0+) - a \in dz) \\ &= W^{(q)}(0+)\nu(a+dz). \end{aligned}$$

For the proof of the second part we should first note that it is known (cf. [6]) that the process X creeps downwards if and only if $\sigma \neq 0$. It thus follows that Y creeps upwards if and only if $\sigma \neq 0$. Henceforth assume that $\sigma \neq 0$. The proof then follows the same reasoning except in (17) one replaces the event $\{a - Y_{\zeta_a-} \in dy, Y_{\zeta_a} - a \in dz\}$ by $\{Y_{\zeta_a} = a\}$ on the left hand side and $\{a - \epsilon(\rho_{a-}) \in dy, \epsilon(\rho_a) - a \in dz\}$ by $\{\epsilon(\rho_a) = a\}$ on the right hand side. Furthermore as a replacement for (18) in the argument we use instead

$$\mathbb{E}(e^{-q\zeta_a}; Y_{\zeta_a} = a) = \frac{\sigma^2}{2} \left\{ \frac{W^{(q)'(a)^2}{W^{(q)}(a)} - W^{(q)''(a)} \right\} \cdot \int_0^\infty \exp \left\{ -\frac{W^{(q)'(a)}}{W^{(q)}(a)} t \right\} dt,$$

which is taken from Theorem 2 of [22]. \square

Proof of Theorem 1.3. We give the proof for the first identity. The proof of the second identity follows along exactly the same lines using the second part of Lemma 2.2 instead and is left as an exercise for the reader.

In a similar spirit to the proof of Lemma 2.2 we may write for a given open interval $B \subset (0, \infty)$,

$$\begin{aligned} & \mathbb{E} \left(e^{-\alpha\kappa - \beta(\tau_0^- - \kappa)}; S_{\tau_0^-}^U \in B, U_{\tau_0^-} \in dy, -U_{\tau_0^-} \in dz \right) \\ &= \mathbb{E}_x \left[\sum_{t \geq 0} \mathbb{I}_{\{L_{t^-}^{-1} < \tau_0^-\}} e^{-\alpha L_{t^-}^{-1}} \mathbb{I}_{\{S_{L_{t^-}^{-1}}^U \in B\}} e^{-\beta(\tau_0^- - L_{t^-}^{-1})} \mathbb{I}_{\{S_{\tau_0^-}^U = S_{L_{t^-}^{-1}}^U\}} \mathbb{I}_{\{U_{\tau_0^-} \in dy, -U_{\tau_0^-} \in dz\}} \right]. \end{aligned}$$

Note however that, on account of the fact that

$$S_{L_{t^-}^{-1}}^U = S_{L_t^{-1}}^U = x + t - \int_x^{L_t^{-1}} \gamma(S_u) dS_u = x + t - \int_x^{x+t} \gamma(y) dy = \bar{\gamma}(x+t)$$

we have that $\{S_{L_t^-}^U \in B\} = \{\bar{\gamma}(x+t) \in B\}$. Note also that $L_t^{-1} = \tau_{\bar{\gamma}(x+t)}^+$ and L^{-1} is quasi left continuous. Hence, applying the compensation formula we have

$$\begin{aligned} & \mathbb{E} \left(e^{-\alpha\kappa - \beta(\tau_0^- - \kappa)}; S_{\tau_0^-}^U \in B, U_{\tau_0^-} \in dy, -U_{\tau_0^-} \in dz \right) \\ &= \mathbb{E}_x \left[\int_0^\infty dt \cdot e^{-\alpha\tau_{\bar{\gamma}(x+t)}^+} \mathbb{I}_{\{\bar{\gamma}(x+t) \in B\}} \mathbb{I}_{\{\tau_{\bar{\gamma}(x+t)}^+ < \tau_0^-\}} \right. \\ & \quad \left. \times n(e^{-\beta\rho_{\bar{\gamma}(x+t)}}; \bar{\gamma}(x+t) - \epsilon(\rho_{\bar{\gamma}(x+t)} -) \in dy, \epsilon(\rho_{\bar{\gamma}(x+t)}) - \bar{\gamma}(x+t) \in dz) \right] \\ &= \int_B \frac{d\theta}{1 - \gamma(\bar{\gamma}^{-1}(\theta))} \mathbb{E}_x \left[e^{-\alpha\tau_\theta^+} \mathbb{I}_{\{\tau_\theta^+ < \tau_0^-\}} \right] n(e^{-\beta\rho_\theta}; \theta - \epsilon(\rho_\theta -) \in dy, \epsilon(\rho_\theta) - \theta \in dz) \end{aligned}$$

where in the final equality we have applied a change of variable. Now making use of the first part of Lemma 2.2 and the conclusion of Theorem 1.1 the result follows. \square

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