

# Strong law of large numbers for supercritical super-diffusions

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## 1. $(L, \beta; D)$ -supercritical branching particle diffusion

- Scatter a  $n$  'initial ancestors' scattered in  $D \subseteq \mathbb{R}^d$  at positions  $x_1, \dots, x_n$ . Write  $\nu(x) = \sum_{i=1}^n \delta_{x_i}(\mathrm{d}x)$  for the measure describing the initial state of the system.

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- From each point, issue an  $L$ -diffusion. Here we take

$$L = \frac{1}{2} \nabla \cdot a(x) \nabla + b(x) \cdot \nabla \quad \text{on} \quad D.$$

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- The resulting process is an (atomic) measure-valued **Markov** process  $\{Z_t : t \geq 0\}$  where  $Z_t(dx) = \sum_{i=1}^{N_t} \delta_{x_i(t)}(dx)$ , where  $\{x_1(t), \dots, x_{N_t}(t)\}$  is the spatial configuration of the  $N_t$  particles that are in existence at time  $t$ .

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- We denote its law by  $\mathbb{P}_\nu$ .

## 2. $(L, \beta; D)$ -supercritical branching particle diffusion

- One way to characterise the evolution of the Markov process  $Z$  is to study its transition semi-group through

$$\mathbb{E}_\nu[e^{-\langle f, Z_t \rangle}] = \prod_{i=1}^n v_f(x_i, t)$$

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- We get

$$\frac{\partial}{\partial x} v_f(x, t) = Lv_f(x, t) + \beta(x)[v_f(x, t)^2 - v_f(x, t)] = 0, \quad x \in D, t \geq 0.$$

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- Can generalise this class of Markov processes and talk about measure-valued processes, such that the measure need not be atomic-valued.

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- We will work with the definition of a superdiffusion on  $D \subseteq \mathbb{R}^d$ ,  $X = \{X_t : t \geq 0\}$  as a Markov process valued in the space of finite measures on  $D$ , denoted by  $\mathcal{M}_F(D)$ , with probabilities  $\{\mathbf{P}_\mu : \mu \in \mathcal{M}_F(D)\}$ , such that

$$\mathbf{E}_\mu[e^{-\langle f, X_t \rangle}] = \exp \left\{ \int_D u_f(x, t) \mu(dx) \right\},$$

where

$$\frac{\partial}{\partial t} u_f(x, t) = Lu_f(x, t) - \psi(u_f(x, t), x), \quad x \in D, t \geq 0$$

with  $u_f(x, 0) = f(x)$ ,  $x \in D$  and

$$\psi(\lambda, x) = -\beta(x)\lambda + \alpha(x)\lambda^2, \quad \lambda \in \mathbb{R}, x \in D,$$

with  $\alpha, \beta \in C^\eta$  and  $\alpha \geq 0$ .

## 4. Linear semigroups

- For both  $(L, \beta; D)$  branching particle diffusions and  $(L, \beta, \alpha; D)$  superprocesses, the linear operator  $L + \beta$  plays a special role.

$$\mathbb{E}_{\delta_x}[\langle f, Z_t \rangle] = \mathbb{E}_{\delta_x}[\langle f, X_t \rangle] = w_f(x, t),$$

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- Spectral properties of  $L + \beta$  tell us something about spatial growth:

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- *Local extinction* is the event that a given (and it turns out subsequently all) compact domain(s),  $B \subset\subset D$  becomes empty:  $\exists T(\omega) < \infty$  such that  $X_{T+t}(B) = 0 \forall t \geq 0$ . [Concept obviously still OK for  $Z$  as well]

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- Theorem:** (Englander-Pinsky '99, Englander-K '04) Local extinction iff  $\lambda_c \leq 0$ . [Theorem doesn't care if you talk about branching particle diffusions or superprocesses]

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$$W_t^\phi(X) := e^{-\lambda_c t} \langle \phi, X_t \rangle, \quad t \geq 0$$

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- Change of measure and spine decomposition (see blackboard). For  $\mu \in \mathcal{M}_F(D)$  such that  $\langle \phi, \mu \rangle < \infty$ ,

$$\frac{d\mathbf{P}_\mu^\phi}{d\mathbf{P}_\mu} \Big|_{\sigma(X_s : s \leq t)} = e^{-\lambda_c t} \frac{\langle \phi, X_t \rangle}{\langle \phi, \mu \rangle}$$

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- It turns out that the spine is a diffusion with generator  $(L + \beta - \lambda_c)^\phi$ : here we use the usual notation for Doob  $h$ -transform to a generator  $A$  (with potential term)

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- $\tilde{\phi}$  is the groundstate of the adjoint of  $L + \beta - \lambda_c$  and the assumption  $\langle \tilde{\phi}, \phi \rangle < \infty$  (and hence  $\langle \tilde{\phi}, \phi \rangle = 1$ ) ensures that the spine is an ergodic diffusion with stationary distribution density  $\tilde{\phi}\phi$ .

## 6. Laws of large numbers

- We also see the the spine by studying the linear semi-group, for “nice”  $f$ ,

$$e^{-\lambda_c t} \mathbf{E}_\mu[\langle f, X_t \rangle] = \int_D \frac{f(y)}{\phi(y)} p^{(L+\beta-\lambda_c)\phi}(x, dy, t) \phi(x) \mu(dx) \xrightarrow{t \rightarrow \infty} \langle f, \tilde{\phi} \rangle \langle \phi, \mu \rangle$$

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- Assuming  $\lambda_c > 0$  and  $\langle \tilde{\phi}, \phi \rangle = 1$  the above limit as well as the fact that  $W_\infty^\phi(X)$  is an UI limit is strongly suggestive that  $\lambda_c$  is in fact the growth rate on compacta in the sense that a limit for

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- Theorem:** Suppose that  $\lambda_c > 0$ ,  $\langle \tilde{\phi}, \phi \rangle = 1$ ,  $\|\alpha\phi\|_\infty < \infty$  and **(Mystery Hypothesis)**, then, for all  $0 \leq f \leq \phi$  and  $\mu \in \mathcal{M}_F(D)$  such that  $\langle \phi, \mu \rangle < \infty$  and  $\mu \in \mathcal{M}_F(D)$ ,

$$\lim_{t \rightarrow \infty} e^{-\lambda_c t} \langle f, X_t \rangle = \langle f, \tilde{\phi} \rangle W_\infty^\phi(X) \quad \mathbf{P}_\mu\text{-a.s.}$$



## 7. The skeleton

- The event  $\mathcal{E} := \{\exists t \geq 0 : X_t(D) = 0\}$  generates the rate function  $\omega$  in the following sense:

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- Moreover, the ground state of this  $(L_0^\omega, \alpha\omega; D)$  branching diffusions is precisely  $\varphi/\omega$  with eigenvalue  $\lambda_c$  because

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- One similarly shows that the ground state of the adjoint of  $L_0^\omega + \alpha\omega$  is  $\omega\tilde{\phi}$ .

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- Remarkably one can prove that  $W_\infty^\phi(X) = W_\infty^{\phi/\omega}(Z)$  almost surely.
- A SLLN for the skeleton would thus read: for  $0 \leq f \leq \phi/\omega$

$$\lim_{t \rightarrow \infty} e^{-\lambda_c t} \langle f, Z_t \rangle = \langle f, \omega\tilde{\phi} \rangle W_\infty^{\phi/\omega}(Z) \quad (*)$$


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## 8. The mystery condition

- As a consequence  $W_t^{\phi/\omega}(Z) := e^{-\lambda_c t} \langle \phi/\omega, Z_t \rangle$ ,  $t \geq 0$ , is a martingale.
- If one changes measure on the skeleton process using the martingale  $W_t^{\phi/\omega}(Z)$ , then a spine decomposition emerges (see blackboard) for which the spine has diffusion operator

$$(L_0^\omega + \alpha\phi - \lambda_c)^{\phi/\omega} = (L + \beta - \lambda_c)^\phi$$

and has stationary distribution  $(\phi/\omega)(\omega\tilde{\phi}) = \phi\tilde{\phi}$ .

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- **Theorem:** Suppose that  $\lambda_c > 0$ ,  $\langle \tilde{\phi}, \phi \rangle = 1$ ,  $\|\alpha\phi\|_\infty < \infty$  and **(\*) holds along all lattice sequences  $\delta\mathbb{N}$ ,  $\delta > 0$** , then, for all  $0 \leq f \leq \phi$  and  $\mu \in \mathcal{M}_F(D)$  such that  $\langle \phi, \mu \rangle < \infty$  and  $\mu \in \mathcal{M}_F(D)$ ,

$$\lim_{t \rightarrow \infty} e^{-\lambda_c t} \langle f, X_t \rangle = \langle f, \tilde{\phi} \rangle W_\infty^\phi(X) \quad \mathbf{P}_\mu\text{-a.s.}$$

## 9. Why the skeleton is a natural approach

- First note, it suffices to prove that for  $0 \leq f \leq \phi$

$$\liminf_{t \rightarrow \infty} e^{-\lambda_c t} \langle f, X_t \rangle \geq \langle f, \tilde{\phi} \rangle W_\infty^\phi(X)$$

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- (See blackboard)

## 10. Examples

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- **Example 1.** [Super-outward-OU process with constant branching] Suppose  $D = \mathbb{R}^d$ ,

$$L = \frac{1}{2}\Delta + \gamma x \cdot \nabla,$$

$\beta$  is a constant valued in  $(\gamma d, \infty)$  and  $\alpha$  is uniformly bounded. Then,

$$\lambda_c = \beta - \gamma d, \quad \phi(x) = (\gamma/\pi)^{d/2} \exp\{-\|x\|^2\}, \quad \tilde{\phi}(x) = 1$$

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. All conditions, in particular (mystery condition), is automatically satisfied.

- **Example 2.** (Continuing unfinished work of Fleischmann & Swart '03). [Super-Fisher-Wright diffusion] Suppose  $D = (0, 1)$ ,  $\beta > 1$  (constant) and  $\alpha(x)$  uniformly bounded and

$$L = \frac{1}{2}x(1-x)\frac{d^2}{dx^2}$$

in which case

$$\lambda_c = \beta - 1, \quad \phi(x) = 6x(1-x), \quad \tilde{\phi}(x) = 1$$