

Skorokhod Embeddings & LPs

Defⁿ 1. Let $(X_t)_{t \geq 0}$ be a stochastic process on \mathbb{R} ,
 and suppose μ is a measure on \mathbb{R} .
 We say a stop. time T is a solⁿ to
 the Skorokhod Embedding Problem (SEP)
 (for μ) if

$$X_T \sim \mu.$$

§1. SEP for BM.

Let B_t a BM, $B_0 = 0$, $F(x) = \mu((-\infty, x])$. Then

$$T := \inf \{t \geq 1 : B_t = F^{-1}(\psi(B_t))\}$$

is an embedding since BM recurrent &

$$P(F^{-1}(\psi(B_t)) \leq x) = P(B_t \leq F^{-1}(F(x)))$$

$$= F(F^{-1}(F(x))) = F(x) = \mu((-\infty, x])$$

However (unless $\mu \sim N(0,1)$) : $E T = \infty$
 (since this is the case for eg. $\mu = \inf \{s : B_s = 1\}$)

Since $B_t^2 - t$ is a m.g., would like $E B_T^2 = E T$.
 i.e. $E T = \int x^2 d\mu$ for "small" stopping times.

Defⁿ 2: ~~A [Monroe]~~ An SE is minimal if
 \forall stop. times S ,

$$S \leq T \text{ a.s.} \text{ \& } B_S \sim \mu$$

$$\Rightarrow S = T \text{ a.s.}$$

Thm [Monroe] If μ centred, ~~any~~ T is minimal
 iff $(B_{t \wedge T})_{t \geq 0}$ is UI.

Ans: find minimal solutions to SEP

For convenience, see ρ is symmetric, integrable with true density & no atoms. Write $\bar{\rho}(x) = \rho(x, \infty)$

Thm: [Vallois]

Let
$$\psi(x) = \int_0^x \frac{x}{\rho(s)} \rho(ds)$$

Then $\tau = \inf \{t \geq 0 : |B_t| \geq \psi^{-1}(L_t)\}$ solves the SEP & is minimal

Proof: Let \mathcal{U} be the space of excursions, and for $\varepsilon \in \mathcal{U}$ write

$$\bar{\varepsilon} = \sup_{s \in \partial \mathcal{D}(\varepsilon)} \varepsilon(s), \quad \underline{\varepsilon} = \inf_{s \in \partial \mathcal{D}(\varepsilon)} \varepsilon(s), \quad |\varepsilon| = \bar{\varepsilon} \vee (-\underline{\varepsilon})$$

where $\partial \mathcal{D}(\varepsilon) = \inf \{u \geq 0 : \varepsilon(u) = 0\}$

Then the excursion measure, n on \mathcal{U} of B_n has:

$$n(\bar{\varepsilon} |\varepsilon| \geq x) = \frac{1}{x}$$

{ Pf: $\mathbb{E} L_{H_n} = x$, also for exc. ~~at~~ above x }

Note that n on \mathcal{U} induces a measure $\tilde{n} = \text{Leb} \times n$ on $\mathbb{R}_+ \times \mathcal{U}$
- No. of excursions ~~up~~ over time.

Let $A_t = \{(s, u) : s \leq t, \sup_{r \leq s} |c_r| \geq \psi^{-1}(s)\}$

then we stop ~~the first time~~ at the smallest t for which our excursion PP places a point in A_t . (On local time scale)

Hence:

$$P(L_T > x) = e^{-\tilde{n}(A_x)}$$

But $\tilde{\pi}(A_\lambda) = \int_0^\lambda n(|z| \geq \psi^{-1}(r)) dr$
 $= \int_0^\lambda \frac{1}{\psi^{-1}(r)} dr$

Note ψ increasing, and so ψ^{-1} is also \Rightarrow
 $\{L_T > \lambda\} = \{|B_T| > \psi^{-1}(\lambda)\}$

~~$P(L_T > \lambda) =$~~ $\int_0^\lambda \frac{1}{\psi^{-1}(r)} dr = \int_0^{\psi^{-1}(\lambda)} \frac{1}{s} \psi'(s) ds$
 $= \int_0^{\psi^{-1}(\lambda)} \frac{f(s)}{\bar{p}(s)} ds = [-\log(\bar{p}(s))]_0^{\psi^{-1}(\lambda)}$
 $= -\log(2\bar{p}(\psi^{-1}(\lambda)))$
 $\Rightarrow P(L_T > \lambda) = 2\bar{p}(\psi^{-1}(\lambda))$

so $P(|B_T| > x) = \frac{1}{2} P(L_T > \psi(x)) = \bar{p}(x)$
 as required.

Now check $B_{\tau \wedge T}$ is UI (and hence minimal)

Note: if $H_N = \inf \{t \geq 0 : |B_t| = N\}$, $(|B_t|_{t \wedge H_N} - L_{t \wedge H_N})$ is UI.

Consider a stop. h. S :

$E |B_{S \wedge T}| \mathbb{1}_{\{S \wedge T \leq H_N\}} \xrightarrow{MC} E |B_{S \wedge T}|$
 \wedge
 $E L_{S \wedge T \wedge H_N} \xrightarrow{MC} E L_{S \wedge T}$

and $E L_{S \wedge T} \leq E L_T = \int_0^\infty P(L_T > \lambda) d\lambda = \int_0^\infty 2\bar{p}(\psi^{-1}(\lambda)) d\lambda$
 $\stackrel{\lambda = \psi^{-1}(u)}{=} \int_0^\infty 2u \bar{p}(u) du = 2E(|B_T|)$

Hence $E |B_{SAT}| \leq 2E (B_T)_+ \quad \forall S \subseteq T.$

$\Rightarrow (B_{SAT})$ is UL.

§2. ~~Long~~ Spectrally Negative Lévy processes.

Now suppose X_t is a SNLP with a Gaussian component, which does not drift to ~~infinity~~ $-\infty$

We wish to derive an excursion theory (from 0) for X_t . Note: we can ~~have~~ ~~se~~ now have excursions which begin positive, but end negative (although not vice versa).

Lemma [Pitovnis + Oloj $\S 07$]: If W is the scale fn of X ,

$$n(\{\bar{\varepsilon} \geq \eta\}) = \frac{1}{W(\eta)}.$$

$$n(\{\bar{\varepsilon} \geq \eta \text{ or } \underline{\varepsilon} \leq -\delta\}) = \frac{W(\eta + \delta)}{W(\eta) W(\delta)}$$

$$n(\{\bar{\varepsilon} = 0, \underline{\varepsilon} \leq -\delta\}) = \frac{\varepsilon^2}{2} \frac{W'(\delta)}{W(\delta)}.$$

"Idea":

Roughly, expect

$$P_0(\tau_\eta^+ < \tau_\xi^-) \sim f(\xi) \times n(\{\bar{\varepsilon} \geq \eta\})$$

\Leftrightarrow as $\xi \rightarrow \infty$,

where $f(\xi)$ \rightarrow as $\xi \rightarrow \infty$ is indep of η .

$$\text{But } P_0(\tau_\eta^+ < \tau_\xi^-) = \frac{W(\xi)}{W(\xi + \eta)} \sim \frac{W(\xi)}{W(\eta)}$$

constant will depend on ~~the~~ scaling of local time.

$$\begin{aligned} \text{Then : } P_0(H_\eta < H_{-\delta}) &= \frac{n(\{\bar{\varepsilon} \geq \eta\})}{n(\{\bar{\varepsilon} \geq \eta \text{ or } \underline{\varepsilon} \leq -\delta\})} \\ \Rightarrow n(\{\bar{\varepsilon} \geq \eta \text{ or } \underline{\varepsilon} \leq -\delta\}) &= \frac{n(\{\bar{\varepsilon} \geq \eta\})}{P_0(\dots)} \\ &= \frac{1}{W(\eta)} \bigg/ \frac{W(\delta)}{W(\eta+\delta)} \\ &= \frac{W(\eta+\delta)}{W(\eta)W(\delta)} \end{aligned}$$

$$\begin{aligned} \& n(\{\bar{\varepsilon} = 0, \underline{\varepsilon} \leq -\delta\}) &= \lim_{\eta \downarrow 0} n(\{\bar{\varepsilon} < \eta \text{ or } \underline{\varepsilon} \leq -\delta\}) \\ &= \lim_{\eta \downarrow 0} [n(\{\bar{\varepsilon} \geq \eta \text{ or } \underline{\varepsilon} \leq -\delta\}) - n(\{\bar{\varepsilon} \geq \eta\})] \\ &= \lim_{\eta \downarrow 0} \left[\frac{1}{W(\eta)} \frac{(W(\eta+\delta) - W(\delta))}{W(\delta)} \right] \\ &= \frac{W'(\delta+)}{W'(0+)W(\delta)} = \frac{\sigma^2}{2} \frac{W'(\delta)}{W(\delta)} \end{aligned}$$

Now suppose p is integrable, ~~non-atomic~~
with positive density ~~$p(dx)$~~ $p(x)$

Theorem [Obloj - Pistorius '09]

Define, for $x < 0 < y$,

$$D_p(y) = \int_0^y w(s) \mu(ds), \quad G_p\left(\frac{x}{\sigma}\right) = \int_x^0 \frac{2}{\sigma^2} \frac{w(-s)}{w'(-s)} \mu(ds)$$

$$g(s) = G_p^{-1}(D_p(s)) \ (\leq 0), \quad f(s) = D_p^{-1}(G_p(s)) \ (\geq 0)$$

$$\psi_+(y) = \int_0^y \frac{w(s) \mu(ds)}{(1 + \bar{p}(s) - \bar{p}(g(s)))}$$

$$\psi_-(x) = \int_{-x}^0 \frac{2}{\sigma^2} \frac{w(-s) \mu(ds)}{w'(-s) (1 + \bar{p}(f(s)) - \bar{p}(g(f(s))))}$$

If $D_p(\infty) = G_p(-\infty)$ — i.e. $\int_0^\infty w(s) \mu(ds) = \int_{-\infty}^0 \frac{2}{\sigma^2} \frac{w(-s)}{w'(-s)} \mu(ds)$

Then $T = \inf \{t \geq 0 : X_t = \psi_+^{-1}(L_t) \text{ or } X_t = -\psi_-^{-1}(L_t) \text{ on an initially -ve excursion}\}$

Then T solves the SEP.

$$\phi_+(l) = \psi_+^{-1}(l)$$

Pf: (1) (2) ~~$g(\phi_+(l))$~~ Write ~~$\phi_+(l) = \psi_+^{-1}(l)$~~

Then $g(\phi_+(l)) = -\phi_-(l)$.

Pf: $\Leftrightarrow \psi_-(-g(u)) = \psi_+(u)$

$$\psi_-(-g(u)) = \frac{2}{\sigma^2} \int_{g(u)}^0 \frac{w(-y) \mu(dy)}{w'(-y) (1 + \bar{p}(f(y)) - \bar{p}(y))}$$

$y = g(s)$
 $s = f(y)$

$$= \int_0^u \frac{w(s) \mu(ds)}{(1 + \bar{p}(s) - \bar{p}(g(s)))}$$

$$[\phi_+(u)]' = \frac{1}{\psi_+'(\phi_+(u))} = \frac{(1 + \bar{p}(\phi_+(u)) - \bar{p}(g(\phi_+(u))))}{\omega(\phi_+(u)) e(\phi_+(u))}$$

$$[\phi_-(u)]' = \frac{\sigma^2}{2} \frac{\omega'(\phi_-(u)) (1 + \bar{p}(f(-\phi_-(u))) - \bar{p}(-\phi_-(u)))}{\omega(\phi_-(u)) e(-\phi_-(u))}$$

$$\frac{d}{ds} g'(s) = -\frac{\sigma^2}{2} \frac{\omega(s) \omega'(-g(s))}{\omega(-g(s))} \frac{e(s)}{e(g(s))}$$

& Note :

$$g(\phi_+(u)) = -\phi_-(u)$$

$$\Rightarrow f^{-1}(\phi_+(u)) = -\phi_-(u)$$

$$\Rightarrow \phi_+(u) = f(-\phi_-(u))$$

Hence :

$$\begin{aligned}
 & [1 + \bar{p}(\phi_+(u)) - \bar{p}(g(\phi_+(u)))]' \\
 &= [1 + \bar{p}(\phi_+(u)) - \bar{p}(g(\phi_+(u)))] \\
 & \quad \times \left\{ -\frac{e(\phi_+(u))}{\omega(\phi_+(u)) e(\phi_+(u))} \cdot \frac{1}{\omega(\phi_+(u)) e(\phi_+(u))} \right. \\
 & \quad \left. - \frac{\sigma^2}{2} \frac{\omega(\phi_+(u)) \omega'(-g(\phi_+(u)))}{\omega(-g(\phi_+(u)))} \frac{e(\phi_+(u))}{e(g(\phi_+(u)))} \cdot \frac{1}{\omega(\phi_+(u)) e(\phi_+(u))} \right. \\
 & \quad \left. \times e(g(\phi_+(u))) \right\} \\
 &= [1 + \bar{p}(\phi_+(u)) - \bar{p}(g(\phi_+(u)))] \\
 & \quad \times \left\{ -\frac{1}{\omega(\phi_+(u))} - \frac{\sigma^2}{2} \frac{\omega'(-g(\phi_+(u)))}{\omega(-g(\phi_+(u)))} \right\}
 \end{aligned}$$

$$\Rightarrow \frac{[1 + \bar{p}(\phi_+(u)) - \bar{p}(g(\phi_+(u)))]'}{[1 + \bar{p}(\phi_+(u)) - \bar{p}(g(\phi_+(u)))]} = - \left[n(\bar{\varepsilon} \geq \phi_+(u)) + n(\bar{\varepsilon} = 0, \underline{\varepsilon} \leq -\phi_-(u)) \right]$$

$$\Rightarrow P(L_T > k) = \exp \left(\int_0^k \frac{[\dots]'}{[\dots]} dl \right)$$

$$= 1 + \bar{p}(\phi_+(k)) - \bar{p}(-\phi_-(k)).$$

In particular, for $x > 0$

$$P(X_T > x) = \int_0^\infty P(L_T > l) n(\bar{\varepsilon} > \phi_+(l)) \mathbb{1}\{\phi_+(l) > x\} dl \quad (*)$$

subst. $u = \phi_+(l), \quad du = \frac{(1 + \bar{p}(\dots))}{w(u) \rho(u)} du dl$

$$= \int_0^\infty \frac{1}{w(u)} w(u) \rho(u) \mathbb{1}\{u > x\} du$$

$$= p(x, \infty)$$

similarly for $P(X_T < y), \quad y < 0.$

To deduce minimality, Bertoin - Le Jan $\Rightarrow \mathbb{E}L_s \geq \int_0^\infty w(u) \rho(u) du$
for all s st. $B_s \sim \rho.$

So if $s \leq T, \quad L_s \leq L_T$

But $\mathbb{E}L_T = \int_0^\infty P(L_T > l) dl \stackrel{(*)}{=} \int_0^\infty w(u) \rho(u) du$

So must have $L_s = L_T$ a.s. i.e. stop on same excursion

Spec neg $\Rightarrow X_s \leq X_T$ if S occurs on the exc.
But $X_s \sim X_T \Rightarrow X_s = X_T.$

~~If~~ stop on -ve excursions, either ~~stops~~ can stop before going below $\phi_- \Rightarrow \neq$ by MP or never stop before $\phi_- \Rightarrow X_s = X_T.$