

A new multiscale finite element method for high-contrast elliptic interface problems

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Oberwolfach, August 2008

Main points

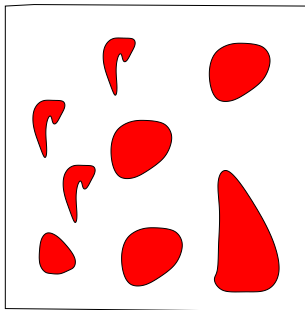
- Elliptic interface problems (**jumping coefficients**)
- **MSFE**: Solve local homogeneous PDEs for basis functions
- **New result**: methods with **optimal convergence** independent of the **contrast** even with “naive meshing”.
- Method involves **new interior boundary conditions** for basis functions.
- Theory involves **new regularity results** for elliptic interface problems
- Method is a **generalisation** of the P1-continuous Galerkin method.
- Theory at present in 2D only.

elliptic interface problems

Find $u \in H_0^1(\Omega)$:

$$\int_{\Omega} \mathcal{A}(x) \nabla u(x) \cdot \nabla v(x) dx = \int_{\Omega} F(x) v(x) dx, \quad v \in H_0^1(\Omega),$$

“High contrast” piecewise constant coefficient \mathcal{A} :



Inclusions: $\Omega_1, \dots, \Omega_m$ $\Omega_0 = \Omega \setminus \cup_{i=1}^m \Omega_i$. Interface Γ .

Scale by $\mathcal{A}_{\min} = \min_x \mathcal{A}(x)$: Find $u \in H_0^1(\Omega)$ such that

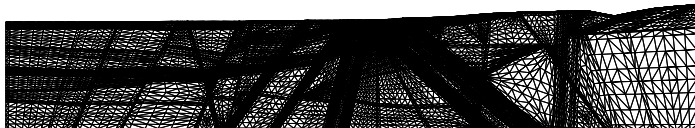
$$a(u, v) := \int_{\Omega} \alpha(x) \nabla u(x) \cdot \nabla v(x) dx = (f, v)_{L_2(\Omega)}, \quad v \in H_0^1(\Omega),$$

with

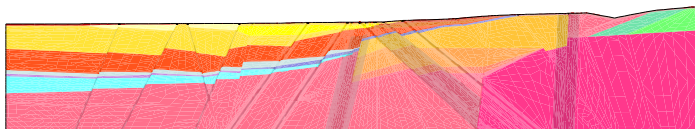
$$\alpha(x) = \frac{1}{\mathcal{A}_{\min}} \mathcal{A}(x), \quad f(x) = \frac{1}{\mathcal{A}_{\min}} F(x).$$

Then $\alpha(x) \geq 1$ and the difficulty is characterised by the **contrast**, a **large parameter**

$$\hat{\alpha} := \frac{\max_x \mathcal{A}(x)}{\min_x \mathcal{A}(x)} \geq 1.$$



- EDZ
- CROWN SPACE
- WASTE VAULTS
- FAULTED GRANITE
- GRANITE
- DEEP SKIDDAW
- N-S SKIDDAW
- DEEP LATTERBARROW
- N-S LATTERBARROW
- FAULTED TOP M-F BVG
- TOP M-F BVG
- FAULTED BLEAWATH BVG
- BLEAWATH BVG
- FAULTED F-H BVG
- F-H BVG
- FAULTED UNDOFF BVG
- UNDOFF BVG
- FAULTED N-S BVG
- N-S BVG
- FAULTED CARB LST
- CARB LST
- FAULTED COLLYHURST
- COLLYHURST
- FAULTED BROCKRAM
- BROCKRAM
- SHALES + EVAP
- FAULTED BNHM
- BOTTOM NHM
- FAULTED DEEP ST BEES
- DEEP ST BEES
- FAULTED N-S ST BEES
- N-S ST BEES
- FAULTED VN-S ST BEES
- VN-S ST BEES
- FAULTED DEEP CALDER
- DEEP CALDER
- FAULTED N-S CALDER
- N-S CALDER
- FAULTED VN-S CALDER
- VN-S CALDER
- MERCIA MUDSTONE
- QUATERNARY



Asymptotic cases

Case I: $\hat{\alpha} := \min_{i=1,\dots,m} \alpha_i \rightarrow \infty$, $\alpha_0 = 1$

Highly permeable inclusions in hardly permeable matrix

Case II: $\hat{\alpha} := \alpha_0 \rightarrow \infty$, $\max_{i=1,\dots,m} \alpha_i \leq \text{Const.}$

Hardly permeable inclusions in highly permeable matrix.

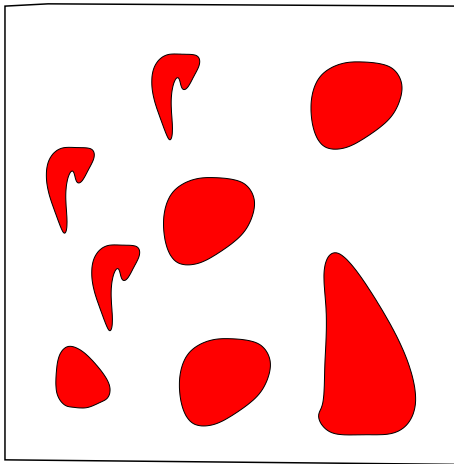
Regularity of solution:

Across an interface Γ separating Ω_- and Ω_+ :

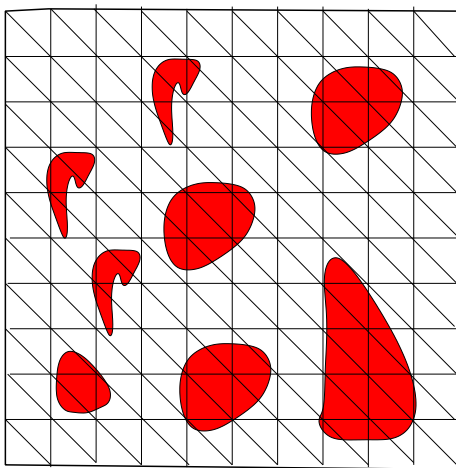
$$\alpha_- \frac{\partial u_-}{\partial n} = \alpha_+ \frac{\partial u_+}{\partial n}$$

Hence $u \in H^{3/2-\epsilon}(\Omega)$. For smooth problems $u \in H^2(\Omega)$

Naive meshing



Naive meshing



Use linear finite elements on the naive mesh.
Finite element space \mathcal{V} .

Naive error estimate Standard linear FEM: seek $u_h \in \mathcal{V}$:

$$a(u_h, v_h) = (f, v_h)_{L^2(\Omega)}, \quad v_h \in \mathcal{V}_h.$$

Energy norm: $|v|_{H^1(\Omega), \alpha} = \left\{ \int_{\Omega} \alpha |\nabla v|^2 \right\}^{1/2}$

$$\begin{aligned} |u - u_h|_{H^1(\Omega), \alpha} &\leq |u - \mathcal{I}_h u|_{H^1(\Omega), \alpha} \\ &\leq C h^{1/2-\epsilon} |u|_{H^{3/2-\epsilon}(\Omega)} = C(\hat{\alpha}) h^{1/2-\epsilon} \end{aligned}$$

(\mathcal{I}_h = piecewise linear interpolant):

Duality argument:

$$\|u - u_h\|_{L_2(\Omega)} \leq C'(\hat{\alpha}) h^{1-2\epsilon}.$$

cf. robust $\mathcal{O}(h)$ and $\mathcal{O}(h^2)$ for smooth problems.

Analysis of $C(\hat{\alpha})$ and $C'(\hat{\alpha})$ requires delicate **regularity theory** (explicit in the **contrast** $\hat{\alpha}$).

“Multiscale” Finite Element Methods

Special finite element space: $\mathcal{V}^{\text{MS}} = \text{span}\{\Phi_p^{\text{MS}}\}$, with

$\Phi_p^{\text{MS}}(x_q) = \delta_{p,q}$ and, for all elements τ :

$$\begin{aligned}\Phi_p^{\text{MS}}|_{\tau} & \text{ is linear} & \tau \cap \Gamma = \emptyset, \\ \Phi_p^{\text{MS}}|_{\tau} & \text{ solves } (*) & \tau \cap \Gamma \neq \emptyset,\end{aligned}$$

Local Homogeneous Problems for the basis functions:

$$\int_{\tau} \alpha \nabla \Phi_p^{\text{MS}} \cdot \nabla v = 0, \quad \text{for all } v \in H_0^1(\tau) \quad (*)$$

Boundary conditions and subgrid **subgrid approximation** needed for (*).

MSFEM: seek $u_h^{\text{MS}} \in \mathcal{V}^{\text{MS}}$:

$$a(u_h^{\text{MS}}, v_h^{\text{MS}}) = (f, v_h^{\text{MS}})_{L^2(\Omega)}, \quad v_h^{\text{MS}} \in \mathcal{V}^{\text{MS}}.$$

The main result

Theorem Assume

- Ω is a convex polygon
- the interface Γ is sufficiently smooth.
- $f \in H^{1/2+\epsilon}(\Omega)$.
- mesh sequence is quasiuniform

Then there exists a choice of boundary condition for Φ_p^{MS} such that

$$(i) \quad |u - u_h^{\text{MS}}|_{H^1(\Omega), \alpha} \lesssim h \left[h|f|_{H^{1/2+\epsilon}(\Omega)}^2 + \|f\|_{L_2(\Omega)}^2 \right]^{1/2},$$

$$(ii) \quad \|u - u_h^{\text{MS}}\|_{L_2(\Omega)} \lesssim h^{2-\epsilon} \left[h|f|_{H^{1/2}(\Omega)}^2 + \|f\|_{L_2(\Omega)}^2 \right]^{1/2}.$$

Hidden constants are independent of h and $\hat{\alpha}$.

requires slightly more regularity on f

non-standard duality argument for the L_2 estimate.

“Subgrid modelling”, e.g. LES in turbulence models, modelling convective storms in NWF, etc..

Hughes et. al. 1995...

Hou and Wu, JCP 1997:

$$-\nabla \cdot a(x/\epsilon) \nabla u = f \quad \text{with } a \text{ periodic, smooth}$$

$$\|u - u_h^{\text{MS}}\|_{H^1(\Omega)} \leq C_1 h \|f\|_{L_2(\Omega)} + C_2 (\epsilon/h)^{1/2} \quad \epsilon \ll h .$$

Many related papers, **Abdulle and E, 03, E & Engquist 04, Efendiev, Hou and Wu, 00, Arbogast & Boyd 06...**

Proofs of accuracy by homogenization arguments but computations on general media

Preconditioning using multiscale coarsening:

IGG, Lechner & Scheichl 07, IGG & Scheichl 07, Pechstein and Scheichl 08. **Proofs of condition number estimates without homogenization structure.**

Solvers and adaptivity (**Without homogenization structure:**)

- Many papers on domain decomposition methods (e.g. **Chan and Mathew, 04**) solvers independent of coefficient jumps (coarse mesh resolves discontinuity).
- Recent proof of robustness of geometric multigrid, **Xu and Zhu, 07**.
- Spectral clustering even without coarse resolution **IGG and Hagger 99, Vuik et. al 00, Aksoylu, IGG, Klie, Scheichl, 08**
- Robustness of a posteriori error estimators, e.g. **Bernardi and Verfürth 00, Ainsworth 05, Vohralik 08**.

Analysis of accuracy of underlying methods ??

Analysis of MSFE: The Main Idea

Optimality:

$$|u - u_h^{\text{MS}}|_{H^1(\Omega), \alpha} \leq |E_h^{\text{MS}}|_{H^1(\Omega), \alpha},$$

MS interpolant $\mathcal{I}_h^{\text{MS}}$

$$E_h^{\text{MS}} := u - \mathcal{I}_h^{\text{MS}}.$$

By **definition** of multiscale basis functions:

$$a_\tau(E_h^{\text{MS}}, v) = a_\tau(u, v) = (f, v)_{L^2(\tau)}, \quad \text{for all } v \in H_0^1(\tau).$$

Simple energy argument:

$$|E_h^{\text{MS}}|_{H^1(\tau), \alpha} \lesssim |\tilde{E}_h^{\text{MS}}|_{H^1(\tau), \alpha} + h_\tau \|f\|_{L_2(\tau)},$$

where $\tilde{E}_h^{\text{MS}} = E_h^{\text{MS}}$ on $\partial\tau$ (**any extension**). Then

$$|E_h^{\text{MS}}|_{H^1(\Omega), \alpha}^2 \lesssim h^2 \left[\sum_\tau h_\tau^{-2} |\tilde{E}_h^{\text{MS}}|_{H^1(\tau), \alpha}^2 + \|f\|_{L_2(\Omega)}^2 \right].$$

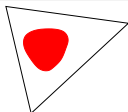
Seek BC s.t. there exists \tilde{E}_h^{MS} yielding

$$\sum_\tau h_\tau^{-2} |\tilde{E}_h^{\text{MS}}|_{H^1(\tau), \alpha}^2 \lesssim h |f|_{H^{1/2+\epsilon}(\Omega)}^2.$$

A simple application: Inclusion inside element

Example: $\hat{\alpha}$ in interior

1 in exterior (Ω_0)



Linear BC's and define $\tilde{E}_h^{\text{MS}} = \begin{cases} E_h^{\text{MS}} & \text{on } \partial\tau \\ 0 & \text{on } \Gamma \end{cases}$.

Trace theorem (scaled):

$$\begin{aligned} |\tilde{E}_h^{\text{MS}}|_{H^1(\tau),\alpha} &\lesssim h_\tau^{-1/2} \|E_h^{\text{MS}}\|_{L^2(\partial\tau)} + h_\tau^{1/2} |E_h^{\text{MS}}|_{H^1(\partial\tau)} \\ &\lesssim h_\tau^{3/2} \|D_t^2 u\|_{L^2(\partial\tau)} \quad \text{tangential derivative} \\ &\lesssim h_\tau^{3/2} \sum_{e \in \mathcal{E}(\tau)} \|D_e^2 u\|_{H^{1/2+\epsilon}(\tau \cap \Omega_0)} \end{aligned}$$

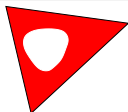
$$\text{So } \sum_\tau h_\tau^{-2} |\tilde{E}_h^{\text{MS}}|_{H^1(\tau),\alpha}^2 \lesssim \sum_\tau h_\tau \sum_{e \in \mathcal{E}(\tau)} \|D_e^2 u\|_{H^{1/2+\epsilon}(\tau \cap \Omega_0)}^2$$

$$\text{(as } \hat{\alpha} \rightarrow \infty) \lesssim h \|f\|_{H^{1/2+\epsilon}(\Omega)}^2 \quad \text{Regularity Theory!!}$$

A simple application: Inclusion inside element

Example: $\hat{\alpha}$ in exterior (Ω_0)

1 in interior



Linear BC's and define $\tilde{E}_h^{\text{MS}} = \begin{cases} E_h^{\text{MS}} & \text{on } \partial\tau \\ 0 & \text{on } \Gamma \end{cases}$.

Trace theorem (scaled):

$$\begin{aligned} |\tilde{E}_h^{\text{MS}}|_{H^1(\tau),\alpha} &\lesssim h_\tau^{-1/2} \hat{\alpha} \|E_h^{\text{MS}}\|_{L^2(\partial\tau)} + h_\tau^{1/2} \hat{\alpha} |E_h^{\text{MS}}|_{H^1(\partial\tau)} \\ &\lesssim h_\tau^{3/2} \hat{\alpha} \|D_t^2 u\|_{L^2(\partial\tau)} \quad \text{tangential derivative} \\ &\lesssim h_\tau^{3/2} \sum_{e \in \mathcal{E}(\tau)} \hat{\alpha} \|D_e^2 u\|_{H^{1/2+\epsilon}(\tau \cap \Omega_0)} \end{aligned}$$

$$\text{So } \sum_{\tau} h_\tau^{-2} |\tilde{E}_h^{\text{MS}}|_{H^1(\tau),\alpha}^2 \lesssim \sum_{\tau} h_\tau \sum_{e \in \mathcal{E}(\tau)} \hat{\alpha} \|D_e^2 u\|_{H^{1/2+\epsilon}(\tau \cap \Omega_0)}^2$$

$$\text{(as } \hat{\alpha} \rightarrow \infty) \lesssim h \|f\|_{H^{1/2+\epsilon}(\Omega)}^2 \quad \text{Regularity Theory!!}$$

Solution flat in regions of high conductivity. (Proof?)

Much more complicated: “cutting through”

Generic case:

Looking for **piecewise linear** boundary condition: Simple Taylor expansion on true solution u :

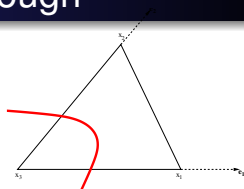
$$r_i^+ \frac{\partial u^+}{\partial e_i}(y_i) + r_i^- \frac{\partial u^-}{\partial e_i}(y_i) = u(x_3) - u(x_1) + \mathcal{O}(h^2) \quad i = 1, 2$$

For gradients \mathbf{g}_i of basis functions on two intersected edges, solve :

$$\begin{aligned} \mathbf{r}_i^T \mathbf{g}_i &= f_i \quad , \quad i = 1, 2 \quad , \\ \mathbf{g}_1 &= B(\theta_1, \theta_2, \beta, \hat{\alpha}) \mathbf{g}_2 \end{aligned}$$

Matrix B depends on angles of intersection, corner and $\hat{\alpha}$

If Γ orthogonal to edges, B reduces to **two independent interface conditions**. (“oscillatory b.c.” of Hou and Wu (97)).



interface cutting through

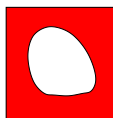
- The recipe leads to **non-conforming** elements, but **averaging** returns conformity without loss of convergence.
- In conforming case $\text{supp}(\Phi_p^{\text{MS}})$ can grow with one extra layer of triangles
- Convergence theorem as before:

$$(i) \quad |u - u_h^{\text{MS}}|_{H^1(\Omega), \alpha} \lesssim h \left[h|f|_{H^{1/2+\epsilon}(\Omega)}^2 + \|f\|_{L_2(\Omega)}^2 \right]^{1/2},$$

$$(ii) \quad \|u - u_h^{\text{MS}}\|_{L_2(\Omega)} \lesssim h^{2-\epsilon} \left[h|f|_{H^{1/2}(\Omega)}^2 + \|f\|_{L_2(\Omega)}^2 \right]^{1/2}.$$

Regularity theory

Smooth well-separated inclusions
Particular case: picture



$$\begin{aligned} -\nabla \cdot \alpha \nabla u &= f && \text{on } \Omega \\ u &= 0 && \text{on } \partial\Omega \end{aligned}$$

Theorem (I.V. Kamotski and V.P. Smyshlyaev)

$$|u|_{H^{2+s}(\Omega_0)} \lesssim \frac{1}{\hat{\alpha}} \|f\|_{H^s(\Omega)} \quad s \geq 0 \quad (1)$$

$$|u|_{H^{2+s}(\Omega_1)} \lesssim \|f\|_{H^s(\Omega)} \quad s \geq 0 \quad (2)$$

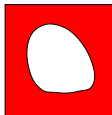
cf. Huang and Zou, J. Differential Equations, 2002

Idea of proof: Introduce \hat{u} solution of

$$-\nabla \cdot \alpha_i \nabla \hat{u} = f_i, \quad \text{on } \Omega_i, \quad i = 0, 1, \quad \hat{u} = 0 \quad \text{on } \partial\Omega, \Gamma$$

decoupled problems, \hat{u} satisfies estimates!

Consider remainder: $\tilde{u} := u - \hat{u}$:



$-\Delta \tilde{u}_i = 0$ on Ω_1 and Ω_0 and $\tilde{u} = 0$ on $\partial\Omega$

Jump condition on interface $\Gamma = \partial\Omega_1$:

$$\hat{\alpha} \frac{\partial \tilde{u}_0}{\partial n} - \frac{\partial \tilde{u}_1}{\partial n} = F \quad := \quad \frac{\partial \hat{u}_1}{\partial n} - \hat{\alpha} \frac{\partial \hat{u}_0}{\partial n} \quad (\dagger)$$

Let $\tilde{v} := \tilde{u}|_{\Gamma}$ and introduce **Dirichlet to Neumann maps** \mathcal{N}_i

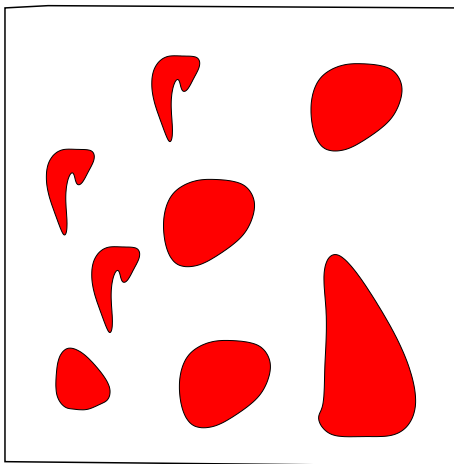
$$\begin{aligned} (\dagger) &\iff (\hat{\alpha}\mathcal{N}_0 - \mathcal{N}_1)\tilde{v} = F \\ &\iff (I - \hat{\alpha}^{-1}\mathcal{N}_0^{-1}\mathcal{N}_1)\tilde{v} = \hat{\alpha}^{-1}\mathcal{N}_0^{-1}F \end{aligned}$$

Contraction mapping ($\hat{\alpha}^{-1} \rightarrow 0$):

$$\begin{aligned} \|\tilde{v}\|_{H^{s+3/2}(\Gamma)} &\lesssim \hat{\alpha}^{-1} \|\mathcal{N}_0^{-1}F\|_{H^{s+3/2}(\Gamma)} \lesssim \hat{\alpha}^{-1} \|F\|_{H^{s+1/2}(\Gamma)} \\ &\lesssim \hat{\alpha}^{-1} \|\hat{u}_1\|_{H^{s+2}(\Omega_1)} + \|\hat{u}_0\|_{H^{s+2}(\Omega_1)} \lesssim \hat{\alpha}^{-1} \|f\|_{H^s(\Omega)} \end{aligned}$$

(In this case $\|\tilde{u}\|_{H^{2+s}(\Omega_0)} = \mathcal{O}(\hat{\alpha}^{-1})$)

Slightly harder case:



Dirichlet to Neumann maps not invertible on “floating” domains.

Numerical Results

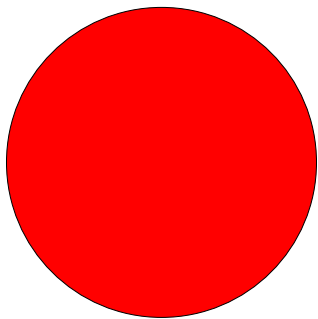
$$\begin{aligned} -\nabla \cdot \alpha \nabla u &= f \quad \text{on } \Omega := [0, 1]^2, \\ u &= g \quad \text{on } \partial\Omega \end{aligned}$$

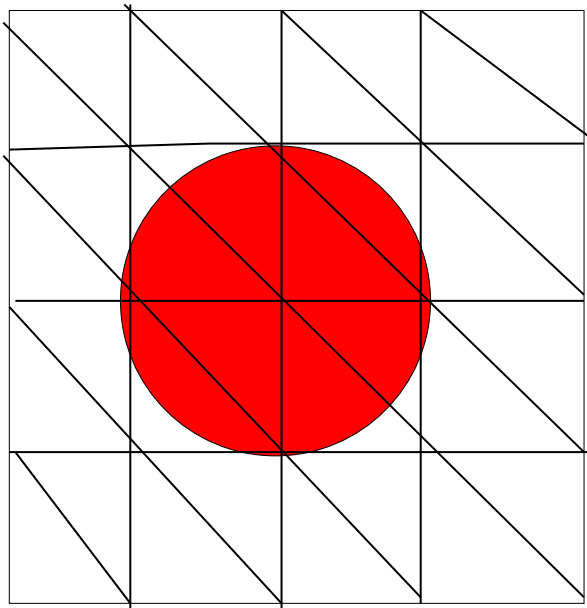
Interface is a circle of radius r_0 ,

$$\alpha(x) = \begin{cases} \alpha_1, & r < r_0 \\ \alpha_0, & r \geq r_0 \end{cases}$$

Exact solution:

$$u(x) = u(r, \theta) = \begin{cases} \frac{r^3}{\alpha_1} & r < r_0 \\ \frac{r^3}{\alpha_0} + \left(\frac{1}{\alpha_1} - \frac{1}{\alpha_0} \right) r_0^3 & r \geq r_0 \end{cases}$$





$$\alpha_1 = 1, \quad \alpha_0 = \hat{\alpha} \rightarrow \infty$$



(Impermeable inclusion in high permeable matrix)

H^1 seminorm errors:

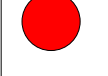
h	$\hat{\alpha} = 10$	$\hat{\alpha} = 10^3$	$\hat{\alpha} = 10^5$
1/8	2.55e-1	2.51e-1	2.54e-1
1/16	1.33e-1	1.24e-1	1.24e-1
1/32	6.22e-2	6.15e-2	6.14e-2
1/64	3.26e-2	3.15e-2	3.07e-2
rate	1.0	1.0	1.0

L_2 errors:

h	$\hat{\alpha} = 10$	$\hat{\alpha} = 10^3$	$\hat{\alpha} = 10^5$
1/8	2.27e-2	2.27e-2	2.29e-2
1/16	5.75e-3	5.76e-3	5.78e-3
1/32	1.45e-3	1.45e-3	1.45e-3
1/64	3.73e-4	3.67e-4	3.63e-4
rate	1.98	1.98	1.99

$$\alpha_0 = 1, \quad \alpha_1 = \hat{\alpha} \rightarrow \infty$$

(Highly permeable inclusion in impermeable matrix)



H^1 seminorm errors:

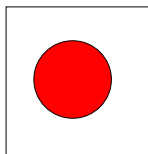
h	$\hat{\alpha} = 10$	$\hat{\alpha} = 10^3$	$\hat{\alpha} = 10^5$
1/8	1.09e-1	5.81e-2	5.90e-2
1/16	4.57e-2	2.75e-2	2.77e-2
1/32	1.43e-2	1.30e-2	1.27e-2
1/64	1.01e-2	6.52e-3	6.10e-3
rate	1.11	1.00	1.09

L_2 errors:

h	$\hat{\alpha} = 10$	$\hat{\alpha} = 10^3$	$\hat{\alpha} = 10^5$
1/8	4.83e-3	3.89e-3	3.89e-3
1/16	1.32e-3	1.10e-3	1.10e-3
1/32	3.32e-4	2.91e-4	2.91e-4
1/64	8.73e-5	7.56e-5	7.53e-5
rate	1.92	1.88	1.88

Solution of subgrid problems

$$h = 1/128, \alpha_1 = \hat{\alpha}, \alpha_0 = 1$$



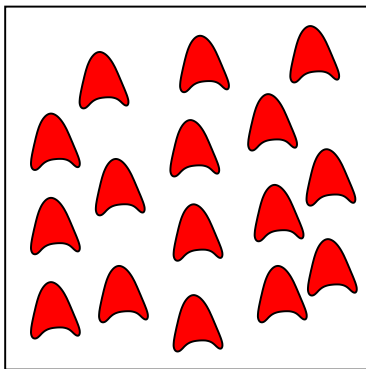
Subgrid problems solved by **Immersed finite element method** (Li, Lin, Wu (2003)).

L^2 errors, where $M = \#$ of subgrid elements

M	$\hat{\alpha} = 10$	$\hat{\alpha} = 10^3$	$\hat{\alpha} = 10^5$
16	1.0307e-4	3.8052e-4	1.4611e-4
64	9.7183e-5	3.8131e-4	7.0831e-5
256	9.6699e-5	3.8188e-4	6.5619e-5
1024	9.6818e-5	3.8216e-4	6.4345e-5

Extensions under construction

Distance between inclusions and distance of inclusions from the boundary are “bad parameters”



With I. Kamotski and V.P. Smyshlyaev (Bath): inclusions separated by $\mathcal{O}(\epsilon)$ and diameter $\mathcal{O}(\epsilon)$.

Conclusion: Summary of results

- Elliptic interface problems with complicated interfaces have **irregular solutions** depending on contrast and interface
- Application of standard FE technology will require **complicated mesh adaptivity** to resolve difficulties
- MSFE can **resolve** these difficulties on **“naive” meshes**.
- The **extra cost** is the solution of subgrid problems on some elements
- Analysis helps explain success of MSFE outside the homogenization framework.
- Regularity theory may help with analysis of **standard methods**.